

Microstructure Mechanics

Dislocation dynamics

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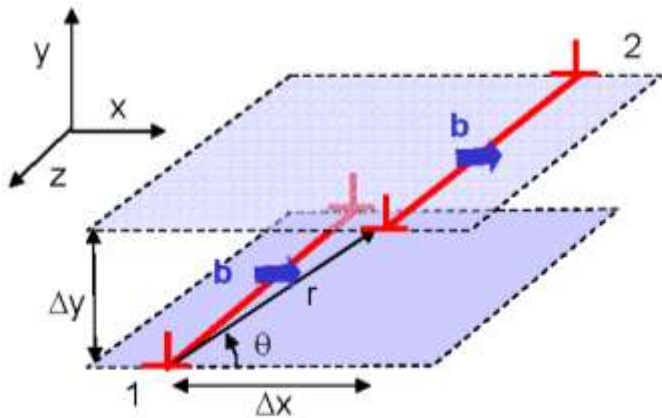


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Dislocation 2 "feels" the stress field of dislocation 1 (and vice versa).

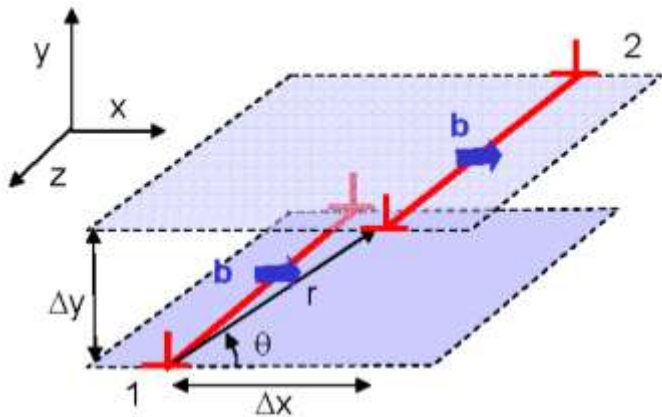
Peach-Koehler Force

$$\vec{F}_{1 \rightarrow 2} = \left(\underline{\underline{\sigma}}^{1 \rightarrow 2} \vec{b}_2 \right) \times \vec{t}_2$$



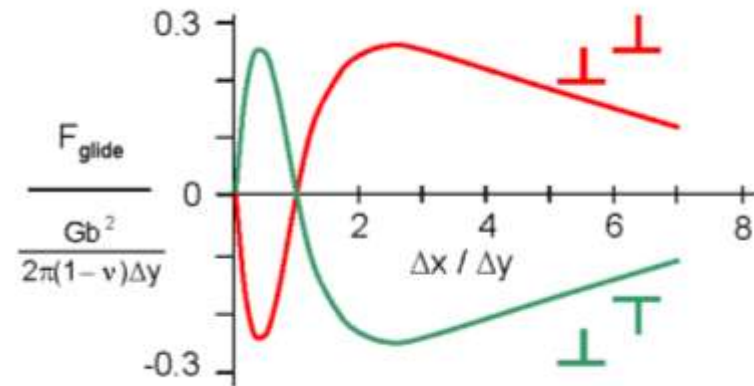
σ_{xy} – produces *glide* force

σ_{xx} – produces *climb* force



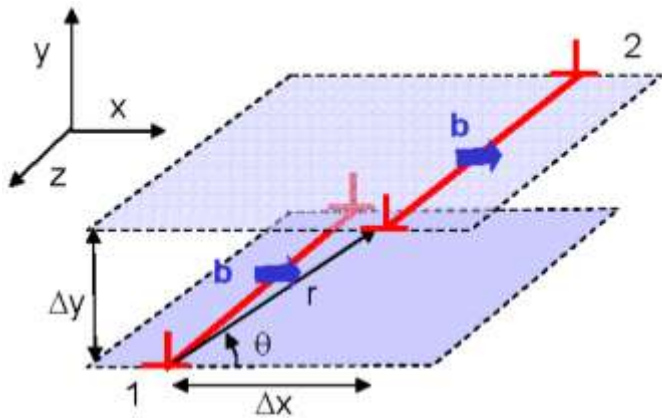
So glide force, resolved onto the slip plane, is:

$$F_{\text{glide}} = \frac{Gb^2}{2\pi(1-\nu)} \frac{\Delta x(\Delta x^2 - \Delta y^2)}{(\Delta x^2 + \Delta y^2)^2}$$



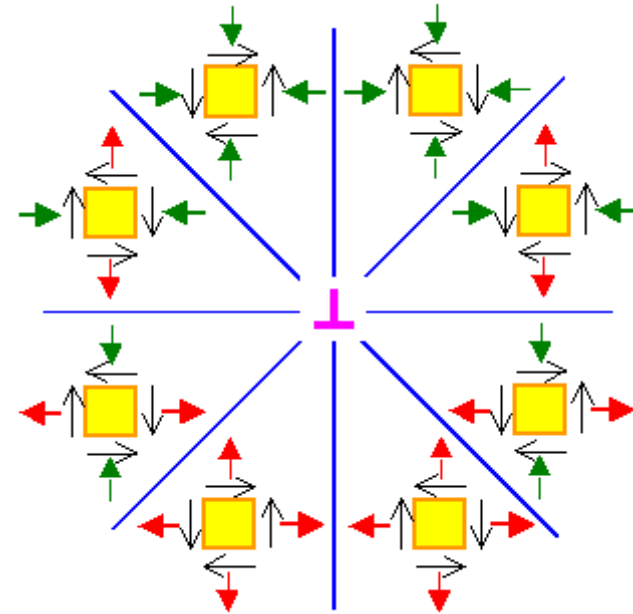
$$\sigma_{xx} = -Dy \frac{3\Delta x^2 + \Delta y^2}{(\Delta x^2 + \Delta y^2)^2}, \quad \text{with: } D = \frac{Gb}{2\pi(1-\nu)}$$

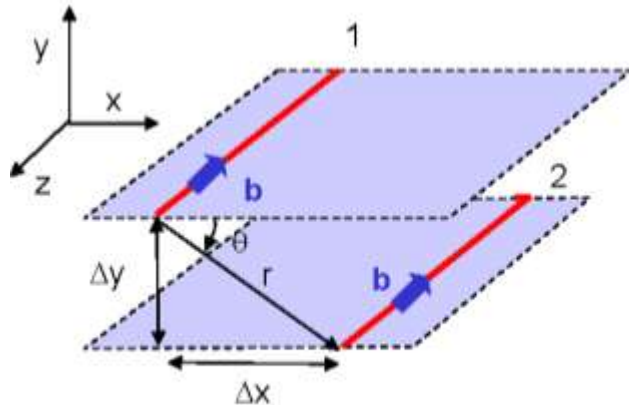
$$\sigma_{xy} = \sigma_{yx} = D\Delta x \frac{\Delta x^2 - \Delta y^2}{(\Delta x^2 + \Delta y^2)^2}$$



$$\sigma_{xx} = -Dy \frac{3\Delta x^2 + \Delta y^2}{(\Delta x^2 + \Delta y^2)^{3/2}}, \quad \text{with: } D = \frac{Gb}{2\pi(1-\nu)}$$

$$\sigma_{xy} = \sigma_{yx} = D\Delta x \frac{\Delta x^2 - \Delta y^2}{(\Delta x^2 + \Delta y^2)^{3/2}}$$





Dislocation 2 "feels" the stress field of dislocation 1 (and vice versa).

$$\sigma_{\theta z} = \sigma_{z\theta} = \frac{Gb}{2\pi r}$$

So force on dislocation 2 from dislocation 1 is:

$$F = \frac{Gb^2}{2\pi r}$$

... but this force acts in the radial direction.

Force on dislocation 2 from dislocation 1, resolved onto the glide plane is:

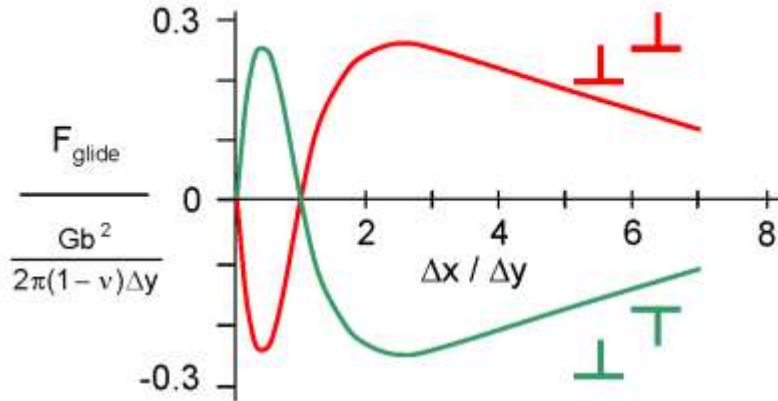
$$F_{\text{res}} = \frac{Gb^2}{2\pi r} \cos \theta$$

Alternatively, we can use the stress field expressed in Cartesian co-ordinates:

$$\begin{aligned} \sigma_{xx} = \sigma_{yy} = \sigma_{zz} = \sigma_{xy} = \sigma_{yx} &= 0 \\ \sigma_{xz} &= -\frac{Gb}{2\pi} \frac{\Delta y}{\Delta x^2 + \Delta y^2} = -\frac{Gb \sin \theta}{2\pi r} \\ \sigma_{yz} &= \frac{Gb}{2\pi} \frac{\Delta x}{\Delta x^2 + \Delta y^2} = \frac{Gb \cos \theta}{2\pi r} \end{aligned}$$

Note that the shear stress acting to shear atoms parallel to **b** above and below the glide plane is σ_{yz} .

$$F_{\text{res}} = \sigma_{yz} b = \frac{Gb^2}{2\pi r} \cos \theta = \frac{Gb^2}{2\pi} \frac{\Delta x}{\Delta x^2 + \Delta y^2}$$



For **like** Burgers vectors:
 $\Delta x = \pm \Delta y$: unstable equilibrium
 $\Delta x = 0$: stable equilibrium

For **opposite** Burgers vectors:
 $\Delta x = \pm \Delta y$: stable equilibrium
 $\Delta x = 0$: unstable equilibrium

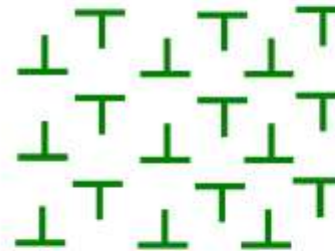
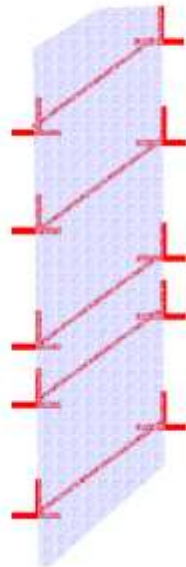
For a set of “**opposite**” Burgers vectors:

There are a large number of possible stable arrangements.

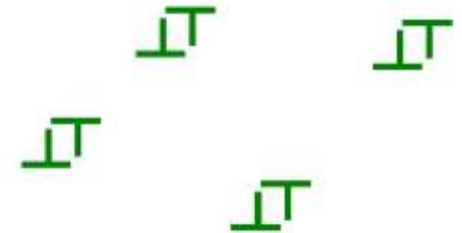
For **like** Burgers vectors:
 Stable array is a planar stack

A low angle tilt boundary.

This arrangement has a strong long-range stress field.



“Taylor lattice”



“Dipole dispersion”

These stable arrangements have minimal *long-range* stress fields.

Discrete Dislocation Dynamics

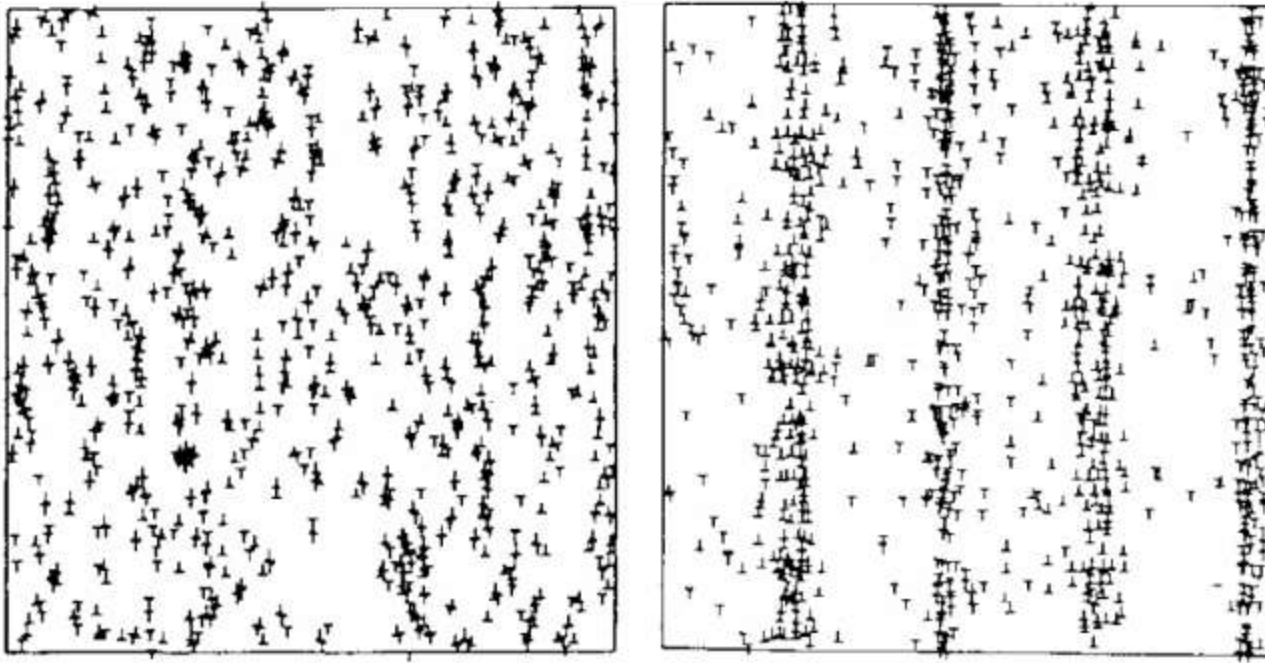
Statistical Dislocation Dynamics



Discrete Dislocation Dynamics

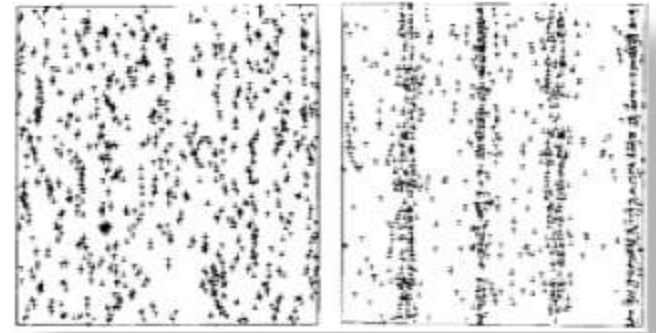
Discrete Dislocation Dynamics in 2D

2D – view parallel to dislocation line



Discrete Dislocation Dynamics in 2D

2D – view parallel to dislocation line



Some questions:

Difference between edge and screw dislocations?

How to do multiplication?

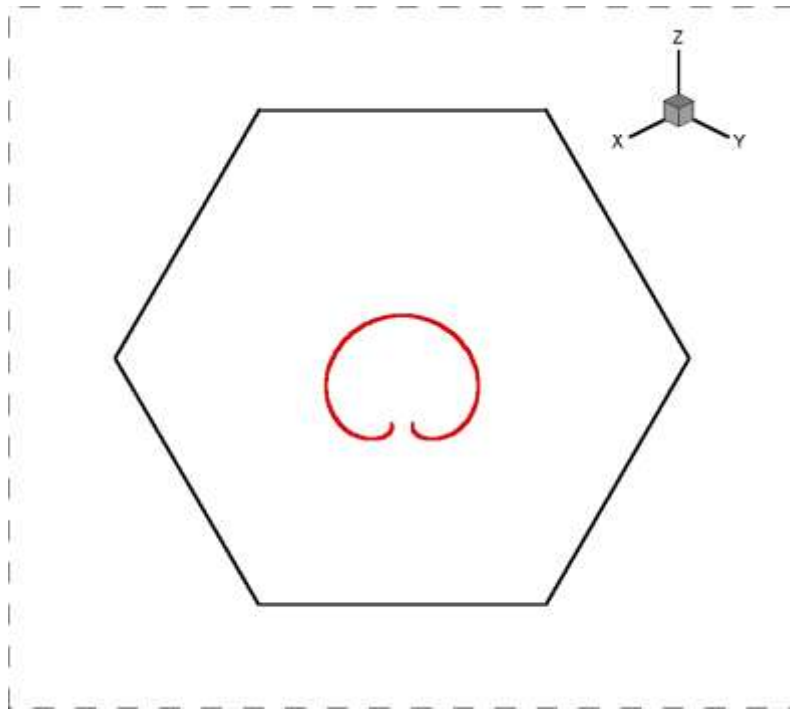
Dislocation bow-out?

Annihilation?

Climbing?

Discrete Dislocation Dynamics in 2D

2D – view into the glide plane





Discrete Dislocation Dynamics in 2D

2D – view into the glide plane

Some questions:

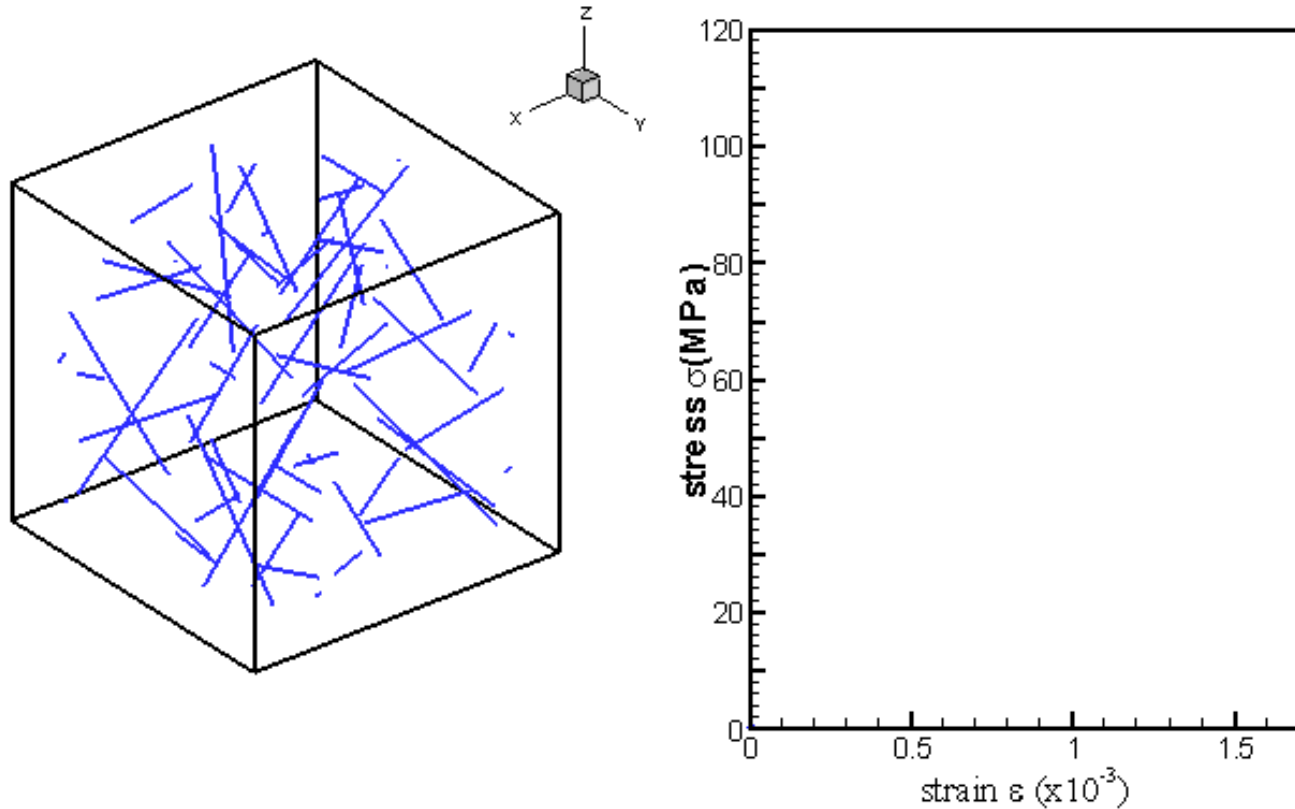
Difference between edge and screw dislocations?

Cross-slip?

Cutting?

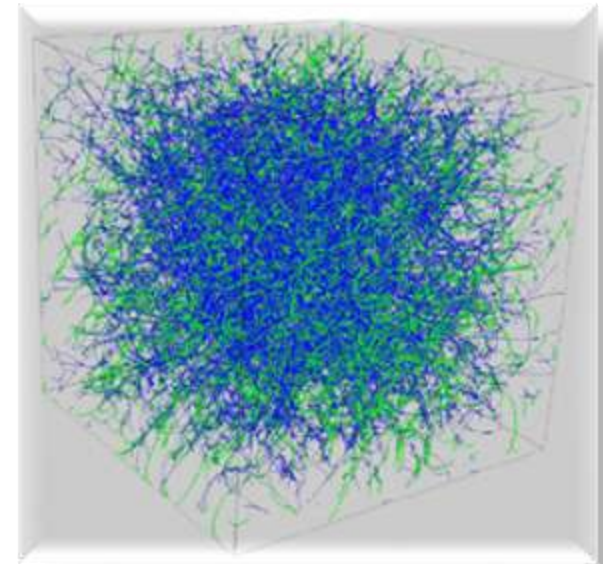
Jog-drag?

3D: DDD (discrete dislocation dynamics)



Discrete Dislocation Dynamics in 3D

Full 3D segment treatment



Some questions:

Difference between edge and screw dislocations?

Junctions?

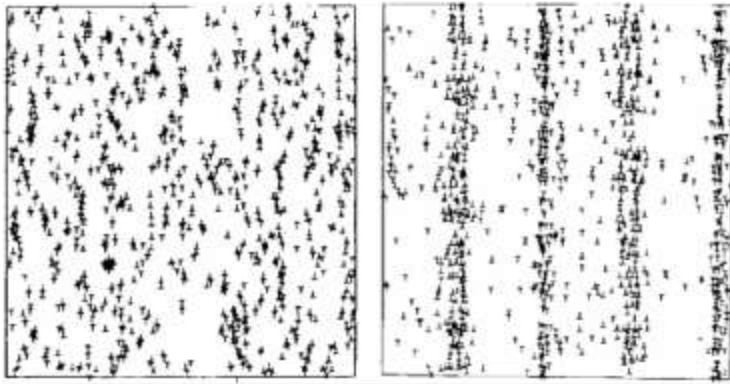
Cutting?

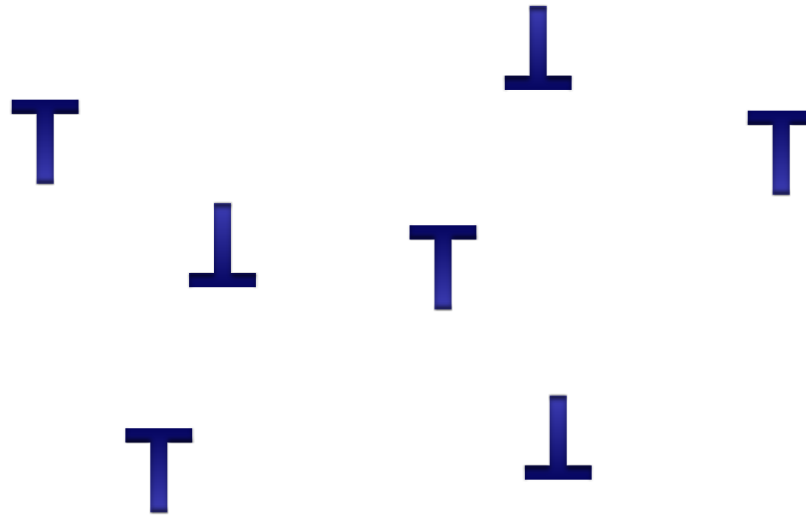
Cores of the dislocations?

Discrete Dislocation Dynamics in 2D

2D – view parallel to dislocation line

Principle procedure





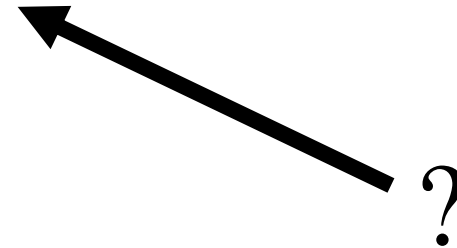
How to proceed?

Stress field of (edge) dislocation

Get coordinates

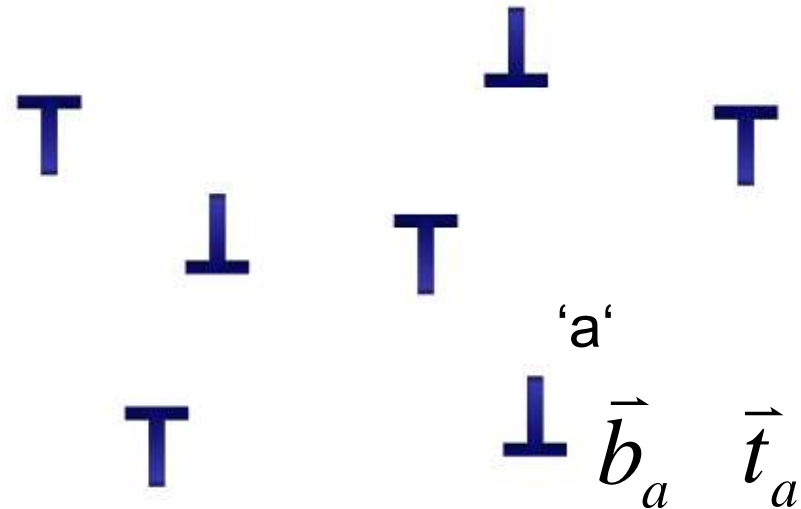
Use Peach Koehler

Move it



Force $\vec{F}_a = \left(\underline{\underline{\sigma}}^{\text{all others} \rightarrow a} \vec{b}_a \right) \times \vec{t}_a$

*Force on dislocation 'a'
by all others*



Force

$$\vec{F}_a = \left(\underline{\underline{\sigma}}^{\text{all others} \rightarrow a} \vec{b}_a \right) \times \vec{t}_a$$

Motion

$$\vec{F} = m \ddot{\vec{x}} + B \dot{\vec{x}} \approx B \dot{\vec{x}}$$

acceleration
friction coefficient (drag)

↓
↓

↑
↑

inertia
velocity

Equilibrium of forces

$$\sum \vec{F}_i = 0$$

$$\sum \vec{F}_i = B\vec{\dot{x}} + \vec{F}_a = 0$$

$$\vec{F}_a = \left(\underline{\underline{\sigma}}^{\text{alle} \rightarrow a} \quad \vec{b}_a \right) \times \vec{t}_a$$

Equilibrium of forces

$$\sum F = 0$$

$$F_{disloc} + F_{self\ force} + F_{extern} + F_{therm} + F_{viscous} + F_{obstacle} + F_{Peierls} + F_{osmotic} + F_{image} + F_{inertia}$$

F_{disloc} : elastic – other dislocations

$F_{obstacle}$: obstacle

$F_{self\ force}$: elastic – self

$F_{Peierls}$: Peierls

F_{extern} : external

$F_{osmotic}$: chemical forces

F_{therm} : Stochastic Langevin

F_{image} : surface forces

$F_{viscous}$: viscous drag

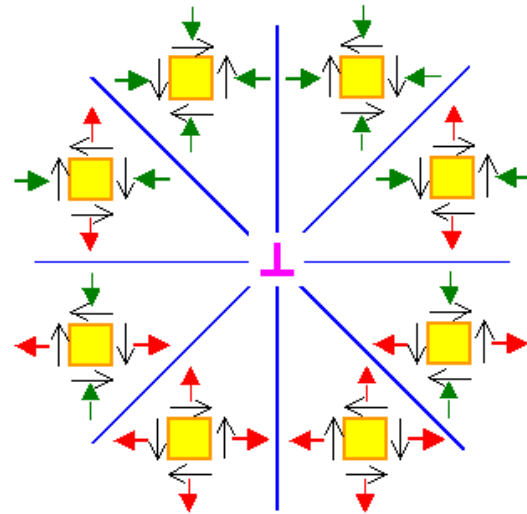
$$\sigma_{xz} = \sigma_{zx} = \sigma_{yz} = \sigma_{zy} = 0$$

$$\sigma_{xx} = -Dy \frac{3x^2 + y^2}{(x^2 + y^2)^2}, \quad \text{with: } D = \frac{Gb}{2\pi(1-\nu)}$$

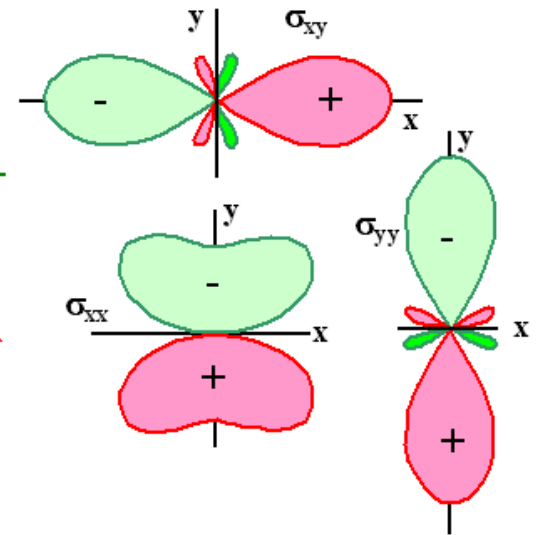
$$\sigma_{yy} = Dy \frac{x^2 - y^2}{(x^2 + y^2)^2}$$

$$\sigma_{xy} = \sigma_{yx} = Dx \frac{x^2 - y^2}{(x^2 + y^2)^2}$$

$$\sigma_{zz} = \nu(\sigma_{xx} + \sigma_{yy})$$



1

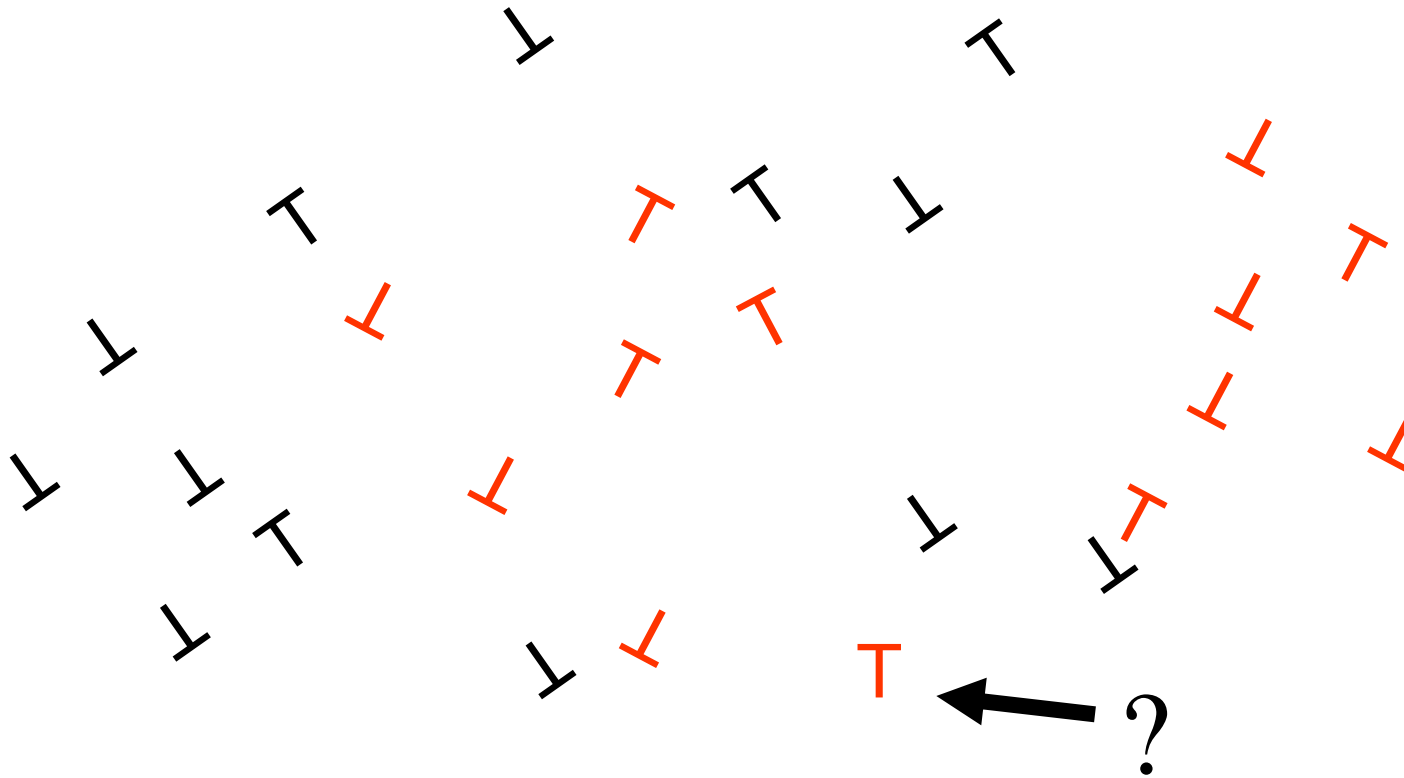


2

$$\vec{F}_a = \left(\underline{\underline{\sigma}}^{\text{all} \rightarrow a} \vec{b}_a \right) \times \vec{t}_a$$

3

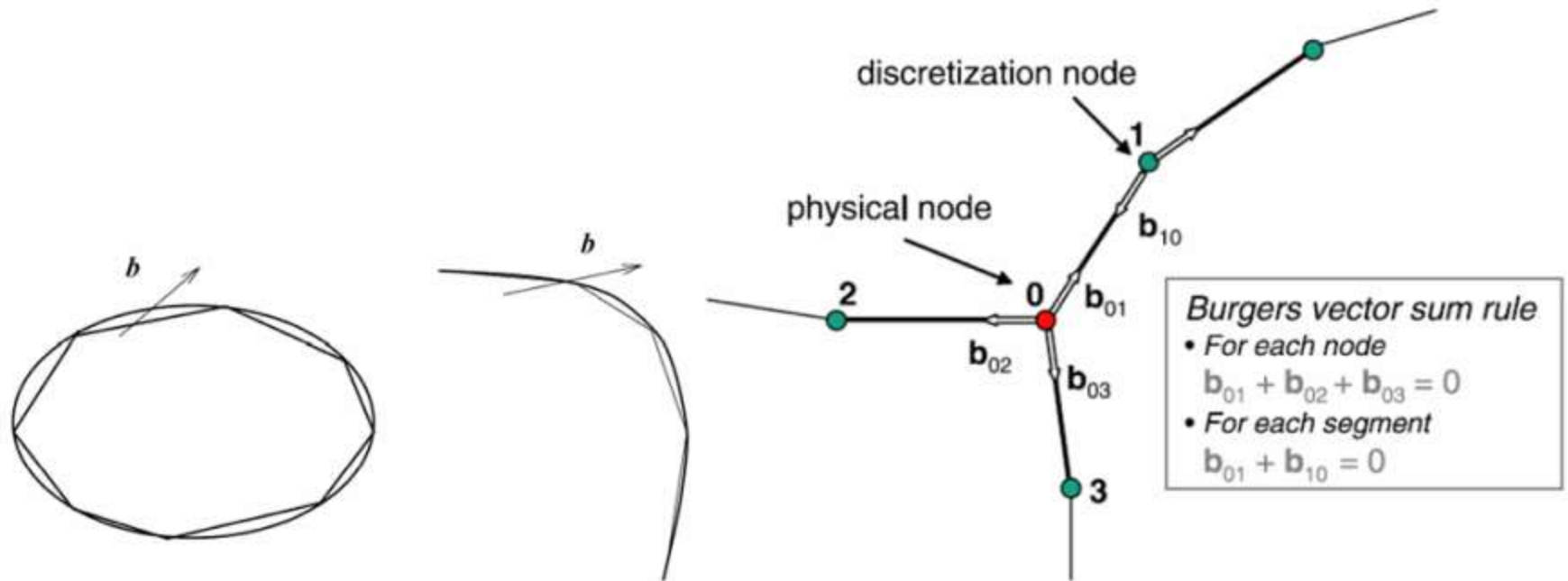
$$\vec{F}_a + B\dot{\vec{x}} + \vec{F}_{\text{external}} = 0$$





- 1) Calculate stress field of machine and of all other dislocations at position of T
- 2) Use Peach-Koehler equation to get force on dislocation
- 3) Integrate with very small time step (explicit) viscous eq. of motion

3D segments and node construction



Annihilation events

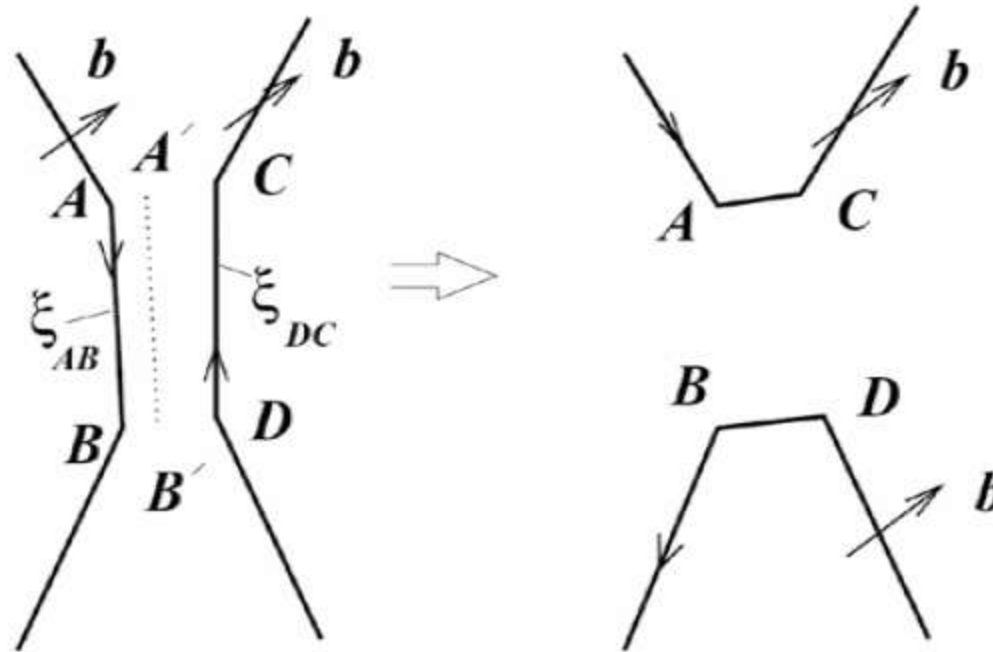


Figure: Annihilation of two attractive dislocations

Jog formation

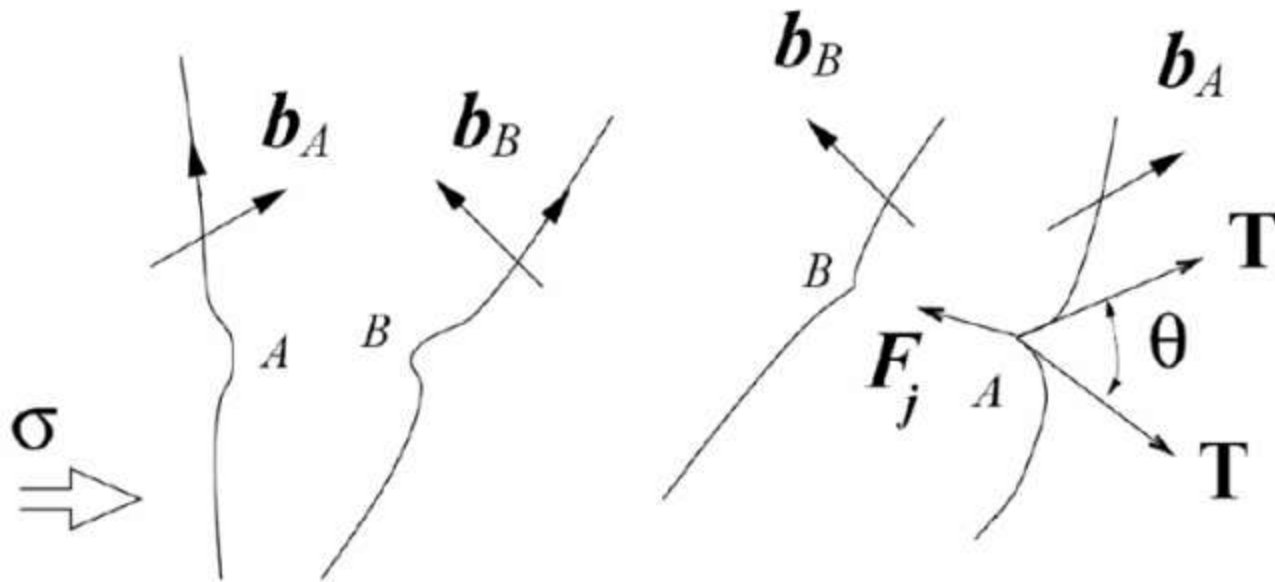
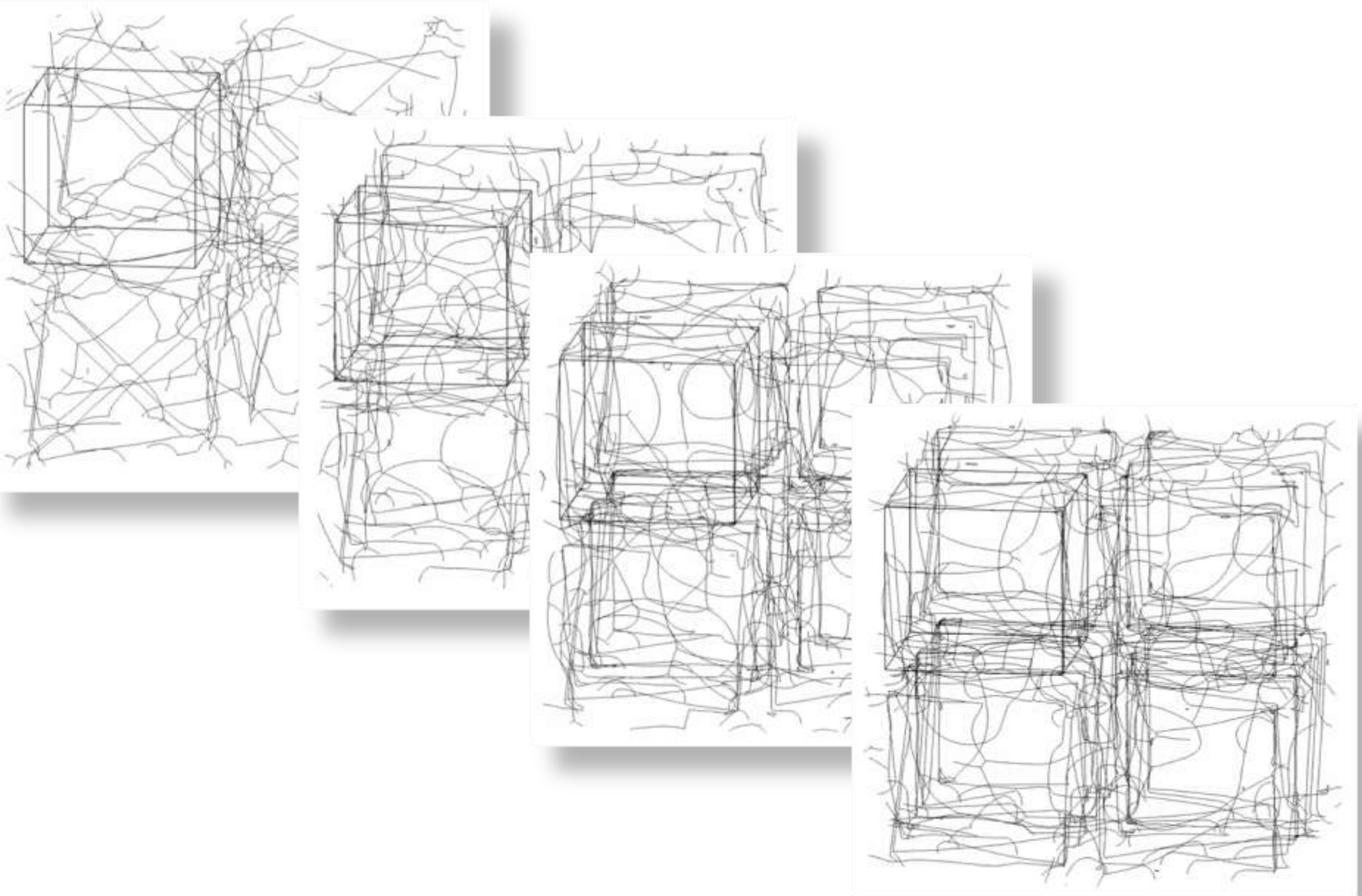
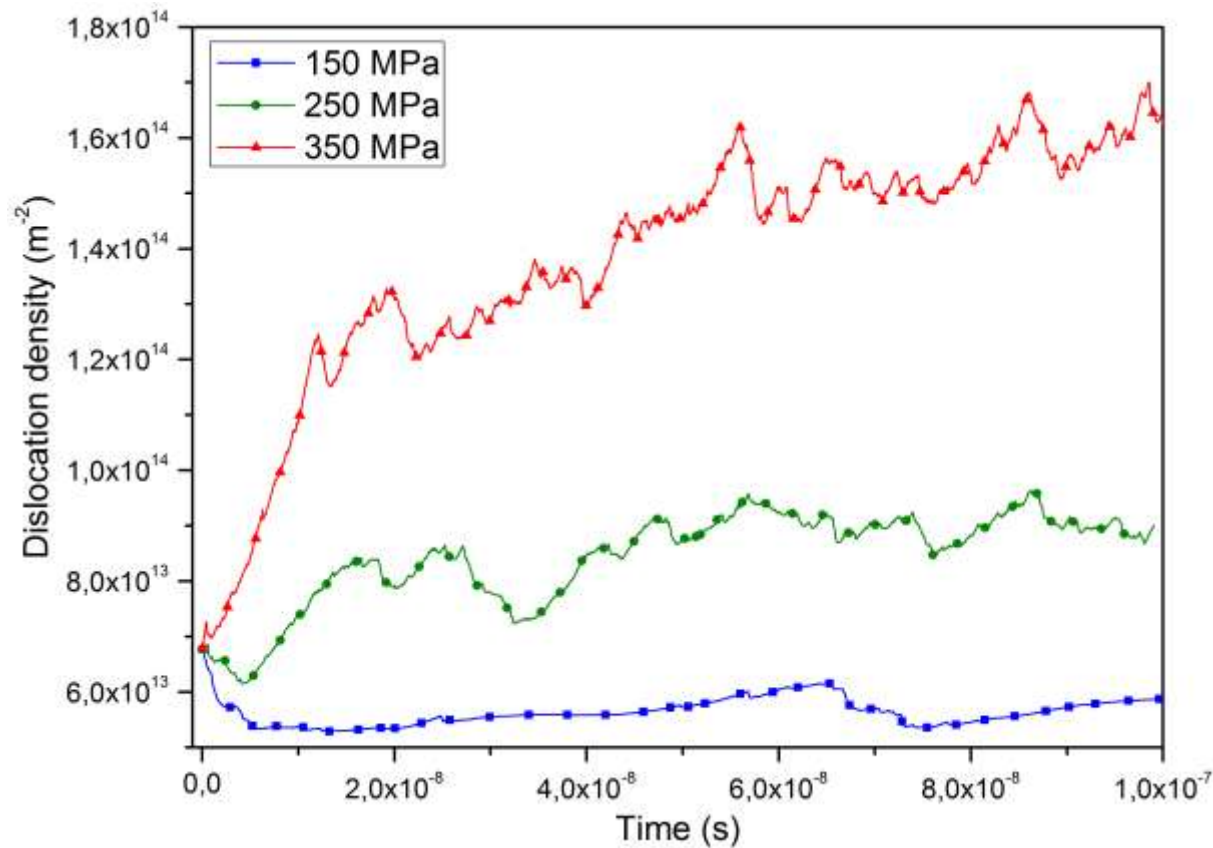


Figure: Jogs are formed when the angle between two attractive dislocations in different planes becomes less than a critical angle θ_c^{jog}

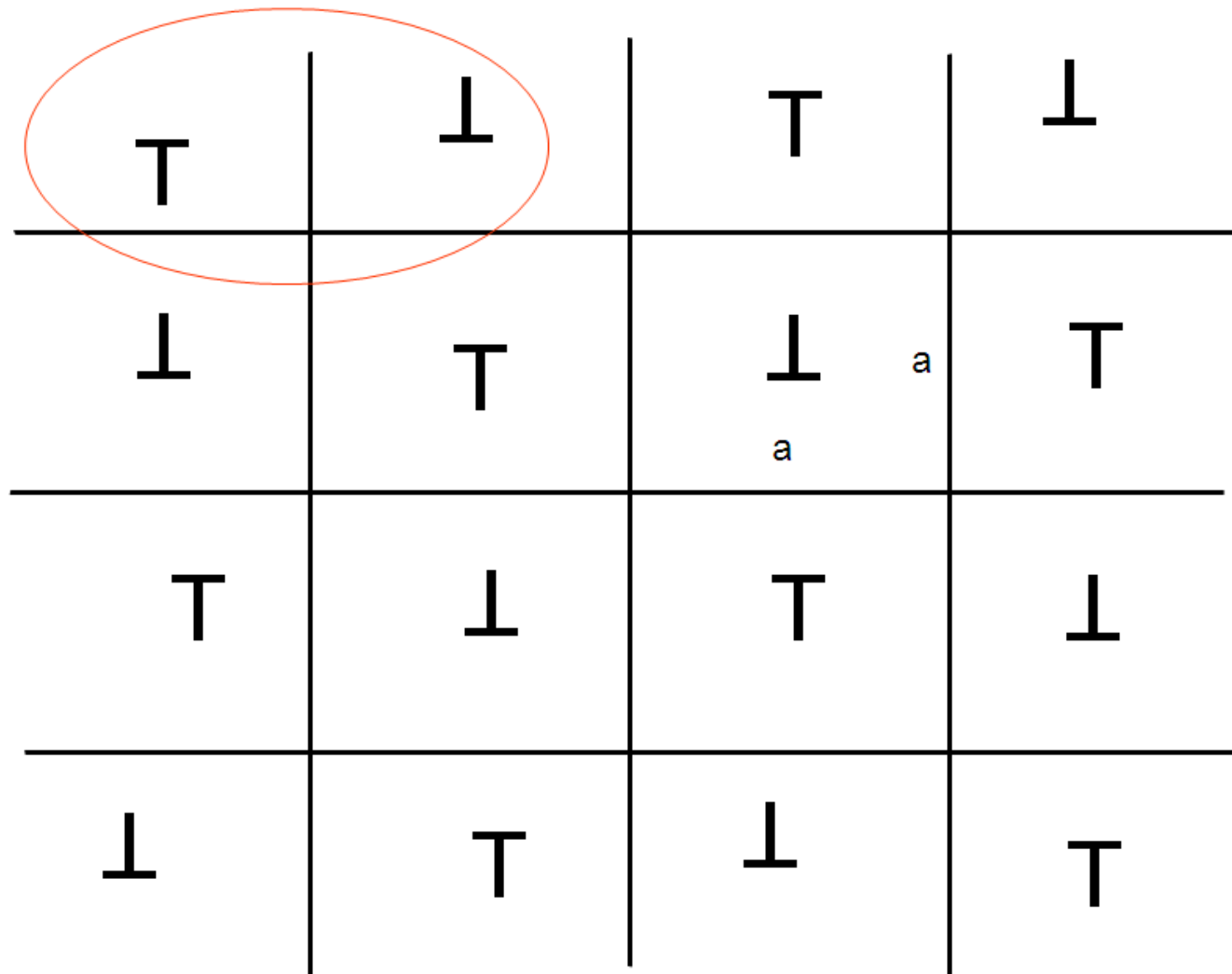


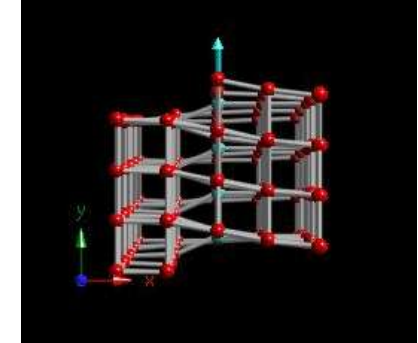
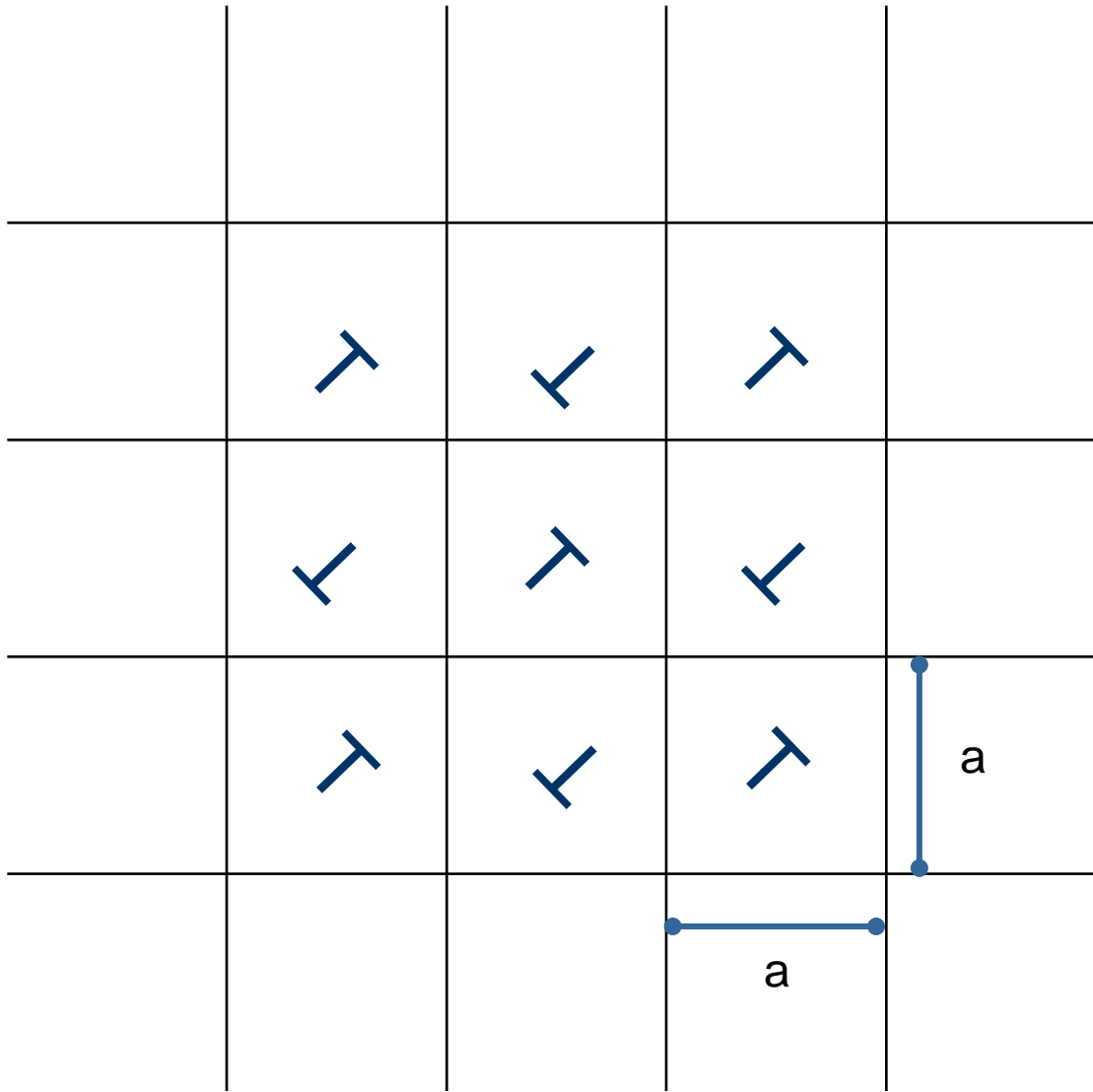




WHY Statistical Dislocation Dynamics ?

- kinetic equation of state
- structure evolution
- coupling to continuum kinematics



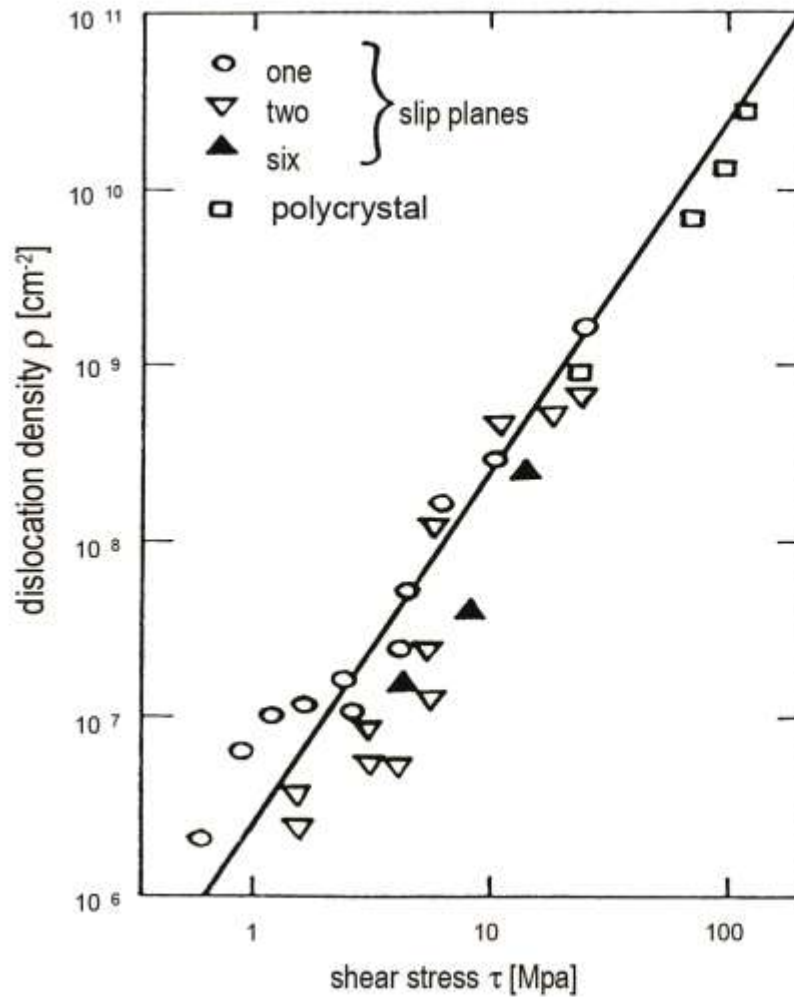


$$\tau = G\gamma = \frac{Gb}{2\pi r}$$

$$\rho = \frac{1}{a^2}$$

$$\tau = \frac{Gb}{2\pi a} = \frac{Gb}{2\pi} \sqrt{\rho}$$

- kinetic equation of state

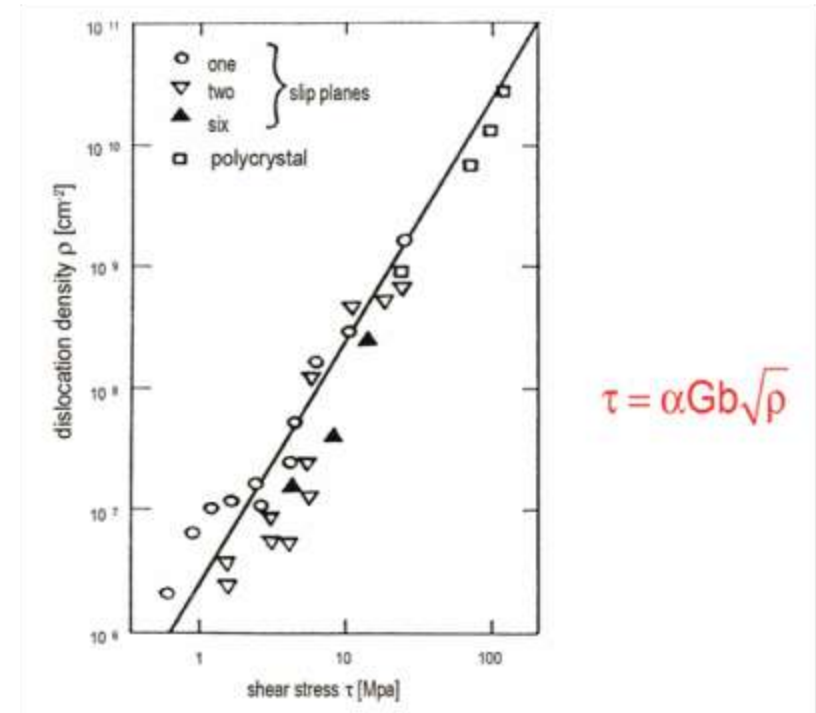


$$\tau = \alpha G b \sqrt{\rho}$$

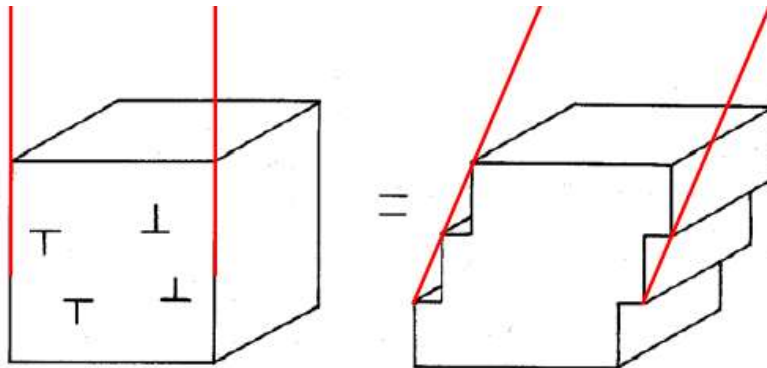
kinetics: collective dislocation behaviour

- kinetic equation of state
- structure evolution

$$\frac{d\rho}{d\gamma} = A\rho^+ + B\rho^-$$



- coupling to imposed shape change



$$\dot{\gamma} = \frac{d\gamma}{dt} = n \frac{dx}{X} \frac{b}{Z} \frac{1}{dt} = \rho_m b v$$